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LETTER TO THE EDITOR

Spiral self-avoiding walks on a triangular lattice: end-to-end distance

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Abstract. An exact solution is obtained for the root-mean-square end-to-end distance R_N of spiral self-avoiding walks with N steps on a triangular lattice. For $N \to \infty$, R_N increases as $(2\pi)^{-1} (6N)^{1/2} \log N$.

Spiral self-avoiding walks (ssaws) on a square lattice were first considered by Privman (1983). It is shown by Blöte and Hilhorst (1984) that, for $N \rightarrow \infty$ the number of N-step spiral walks increases as

$$S_N \sim 2^{-2} 3^{-5/4} \pi N^{-7/4} \exp[2\pi (N/3)^{1/2}],$$
 (1)

and their root-mean-square end-to-end distance behaves as

$$R_N \sim (2\pi)^{-1} (3N)^{1/2} \log N.$$
⁽²⁾

Similar results for S_N are derived independently by Guttmann and Wormald (1984), Joyce (1984), and Guttmann and Hirschhorn (1984). Recently one of us (Lin 1985) considered ssaws on a triangular lattice and it is shown that, for $N \rightarrow \infty$, S_N increases as

$$S_N \sim 2^{1/4} 3^{-7/4} \pi N^{-5/4} \exp[\pi (2N/3)^{1/2}].$$
(3)

In a recent letter, Joyce and Brak (1985) have obtained the complete asymptotic expansion for S_{N} . In this letter we report an exact result for the end-to-end distance of ssaws on a triangular lattice. We shall refer to Lin's paper as I and follow his notations unless specified otherwise.

Consider first the simpler problem of the subclass of ssaws which only spirals outward (see figure 1 of I). The generating function is

$$G^{*}(z) = \sum_{N=1}^{\infty} S_{N}^{*} z^{N}$$

= $\sum_{L=0}^{\infty} \sum_{0 < x_{1} \dots < x_{L}} \sum_{0 < X_{L+1}} z^{x_{1} + \dots + x_{L+1}}$
= $z(1-z)^{-1} \sum_{L=0}^{\infty} g_{L}(z)$ (4)

where

$$g_0(z) \equiv 1$$

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and

$$g_L(z) = \prod_{l=1}^{L} z^l (1-z^l)^{-1}, \qquad L = 1, 2, \dots$$
 (5)

Following exactly the same procedure as Blöte and Hilhorst (1984), it can be shown by the method of steepest descent that for $z = 1 - \varepsilon$ and $\varepsilon \to 0^+$,

$$G^*(z) \sim 2^{-1/2} \varepsilon^{-1} \exp(\pi^2/12\varepsilon).$$
 (6)

For $\varepsilon \rightarrow 0$, the average chain length is

$$N^*(z) = d \log G^*(z) / d \log z \sim \pi^2 / 12\varepsilon^2$$
(7)

and the relative fluctuations satisfy

$$\Delta N^* / N^* \sim (N^*)^{-1/4}.$$
(8)

The squared end-to-end distance is

$$r_L^2 = \frac{1}{2} [(A_L - B_L)^2 + (B_L - C_L)^2 + (C_L - A_L)^2]$$
(9)

where

$$A_L = x_1 + x_4 + \ldots = \sum_{k=0}^{\lfloor L/3 \rfloor} x_{3k+1},$$

$$B_L = x_2 + x_5 + \ldots = \sum_{k=0}^{\lfloor (L-1)/3 \rfloor} x_{3k+2},$$

$$C_L = x_3 + x_6 + \ldots = \sum_{k=1}^{\lfloor (L+1)/3 \rfloor} x_{3k},$$

and [n] denotes the integer part of the real number n. When the expression of r^2 is inserted inside the summations of (4), we have (for $\epsilon \to 0$)

$$F(z) \equiv \sum_{N=1}^{\infty} S_N^* (R_N^*)^2 z^N \sim (R_{N^*}^*)^2 \sum S_N^* z^N$$
(10)

where R_N^* is the root-mean-square end-to-end distance.

We define $m_1 = x_1$, $m_k = x_k - x_{k-1}(k > 1)$ and write

$$2r_{L}^{2} = \left(\sum_{k=0}^{\infty} m_{L-3k}\right)^{2} + \left(x_{L+1} - \sum_{k=0}^{\infty} m_{L-1-3k}\right)^{2} + \left(x_{L+1} - \sum_{k=0}^{\infty} (m_{L-3k} + m_{L-1-3k})\right)^{2}.$$
 (11)

Using the identities

$$\sum_{m=1}^{\infty} m z^m = (1-z)^{-1} \sum_{m=1}^{\infty} z^m,$$
(12)

$$\sum_{m=1}^{\infty} m^2 z^m = (1+z)(1-z)^{-2} \sum_{m=1}^{\infty} z^m,$$
(13)

we have

$$F(z) = z(1-z)^{-1} \sum_{L} R_{L}(z) g_{L}(z)$$
(14)

where R_L is given by

$$R_{L}(z) = \frac{1}{2} \left(\sum_{k} (1 - z^{3k+1})^{-1} \right)^{2} + \frac{1}{2} \left((1 - z)^{-1} - \sum_{k'} (1 - z^{3k'+2})^{-1} \right)^{2} + \frac{1}{2} \left((1 - z)^{-1} - \sum_{k} (1 - z^{3k+1})^{-1} - \sum_{k'} (1 - z^{3k'+2})^{-1} \right)^{2} + z (1 - z)^{-2} + \sum_{k} z^{3k+1} (1 - z^{3k+1})^{-2} + \sum_{k'} z^{3k'+2} (1 - z^{3k'+2})^{-2},$$

$$\sum_{k} = \sum_{k=0}^{\lfloor (L-1)/3 \rfloor} \sum_{k'=0} \sum_{k'=0}^{\lfloor (L-2)/3 \rfloor} .$$
(15)

It has been pointed out by Blöte and Hilhorst (1984) that at fixed z the quantity $g_L(z)$ takes maximum value for $L = L_0(z) \sim \log 2/\epsilon$. Consequently we have

$$F(z) \sim R_{L_0}(z) G^*(z).$$
 (16)

To calculate the leading term in R_{L_0} for $\varepsilon \to 0$, we replace the summations in (15) by integrals and it is elementary to show that

$$R_{L_0}(z) \sim (\log \varepsilon)^2 / 3\varepsilon^2. \tag{17}$$

It follows from (7), (10) and (16) that for $N \rightarrow \infty$

$$R_N^* \sim \pi^{-1} N^{1/2} \log N. \tag{18}$$

We now consider the proglem of all ssaws. It follows from I that, for $z = 1 - \varepsilon$ and $\varepsilon \to 0$, the generating function behaves like

$$G(z) \sim 2G_6(z) - G_1(z) - G_2(z) \tag{19}$$

and

$$2G_1 \sim 6G_2 \sim G_6 \sim \varepsilon \exp(\pi^2/6\varepsilon).$$
⁽²⁰⁾

In analogy to (7), we have

$$N(z) \sim \pi^2/6\varepsilon^2. \tag{21}$$

Consider $G_1(z)$ first (see figure 3 of I). We deote the segment lengths by the integers

$$x_1, x_2, \ldots, x_{L-1}, x_L + y_{L'-1}, y_{L'-1}, \ldots, y_1,$$
 (22)

where

$$0 < x_1 < \ldots < x_L, \qquad 0 < y_1 < \ldots < y_{L'-1}.$$

The length of the longest segment is defined to be $x_L + y_{L'-1}$. We have

$$G_{1}(z) = \sum_{L,L'=1}^{\infty} \sum_{x_{1},...,x_{L}} \sum_{y_{1},...,y_{L'-1}} z^{x_{1}+...+x_{L}+2y_{L'-1}+y_{L-2}+...+y_{1}}$$

$$= z^{-1}(1-z) \sum_{L,L'} g_{L}(z)g_{L}(z)$$

$$= z^{-1}(1-z) \left(\sum_{L=1}^{\infty} g_{L}(z)\right)^{2}.$$
 (23)

The squared end-to-end distance is

$$(r_{L,L'})^2 = \frac{1}{2}[(A-B)^2 + (B-C)^2 + (C-A)^2]$$
(24)

where

$$A = \sum_{k=0}^{\infty} x_{L-1-3k} + \sum_{k'=0}^{\infty} y_{L'-2-3k'}$$

$$B = \sum_{k=0}^{\infty} x_{L-2-3k} + \sum_{k'=0}^{\infty} y_{L'-1-3k'}$$

$$C = \sum_{k=0}^{\infty} x_{L-3k} + y_{L'-1} + \sum_{k'=0}^{\infty} y_{L'-3-3k'}.$$
(25)

We define $m_i = x_i - x_{i-1}$ and $n_j = y_j - y_{j-1}$ $(m_1 = x_1 \text{ and } n_1 = y_1)$ where m_i and n_j are independent integer variables. We have

$$A - B = \sum_{k=0}^{\infty} m_{L-1-3k} - \sum_{k'} n_{L'-1-3k'}$$

$$C - B = \sum_{k=0}^{\infty} (m_{L-3k} + m_{L-1-3k}) + \sum_{k'=0}^{\infty} n_{L'-3-3k'}$$

$$C - A = \sum_{k=0}^{\infty} m_{L-3k} + \sum_{k'=0}^{\infty} (n_{L'-3-3k'} + n_{L'-1-3k'}).$$
(26)

When the expression of r^2 is inserted inside the summations of (23), we have for $\varepsilon \to 0$

$$(\mathbf{R}_N)^2 \sim (\log \varepsilon)^2 / \varepsilon^2 \tag{27}$$

where $N(\varepsilon)$ is the average chain length given by (21). Considerations of $G_2(z)$ and $G_6(z)$ yield identical results. Therefore finally we have

$$R_N \sim (2\pi)^{-1} (6N)^{1/2} \log N \tag{28}$$

for $N \to \infty$.

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References

Blöte H W J and Hilhorst H J 1984 J. Phys. A: Math. Gen. 17 L111-5 Guttmann A J and Hirschhorn M 1984 J. Phys. A: Math. Gen. 17 3613-4 Guttmann A J and Wormald N C 1984 J. Phys. A: Math. Gen. 17 L271-4 Joyce G S 1984 J. Phys. A: Math. Gen. 17 L463-7 Joyce G S and Brak R 1985 J. Phys. A: Math. Gen. 18 L293-8 Lin K Y 1985 J. Phys. A: Math. Gen. 18 L145-8 Privman V 1983 J. Phys. A: Math. Gen. 16 L571-3